Electroweak Standard Model

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02. - 12. September 2003

- Lecture 1: Formulation of the Electroweak Standard Model
- Lecture 2: Renormalization
- Lecture 3: Radiative Corrections in the Electroweak Standard Model

Lecture 1

Formulation of Electroweak Standard Model

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- Electroweak phenomenology before the Standard Model
- Basic Principles of the Standard Model
  - gauge invariance
  - spontaneous symmetry breaking
- Lagrangian of the Electroweak Standard Model

Some phenomenological facts

- discovery of radioactivity (Becquerel 1896)
  β-decay of heavy nuclei: $\mathrm{n} \to \mathrm{p} + \mathrm{e}^- + \bar{\nu}_\mathrm{e}$,
  $\mathrm{p} \to \mathrm{n} + \mathrm{e}^+ + \nu_\mathrm{e}$ (not possible for free protons)
- terminology “weak”: interaction at low energy has very short range
  $\hookrightarrow$ long life time of weakly decaying particles:

strong int.:	$\rho \to 2\pi$,	$\tau \sim 10^{-22}$s
elmg. int.:	$\pi \to 2\gamma$,	$\tau \sim 10^{-16}$s
weak int.:	$\pi^- \to \mu^- + \bar{\nu}_\mu$	$\tau \sim 10^{-8}$s
	$\mu^- \to \mathrm{e}^- + \bar{\nu}_\mathrm{e} + \nu_\mu$,	$\tau \sim 10^{-6}$s

- lepton-number conservation: $\mu^- \not\to \mathrm{e}^- + \gamma$ (BR $\lesssim 10^{-11}$)
- parity violation (predicted by Lee, Yang 1956, detected by Wu 1957)
  e.g.: $\pi^+ \to \mu^+ + \nu_\mu$
  μ^+ always left-handed
  $^{60}\mathrm{Co} \to {}^{60}\mathrm{Ni}^* + \mathrm{e}^- + \bar{\nu}_\mathrm{e}$
  electrons are emitted predominantly
  in direction of spin of $^{60}\mathrm{Co}$
- CP violation (Cronin, Fitch 1964)

$$\begin{array}{lcc} & K_\mathrm{L} \quad \to & 2\pi \\ \mathrm{CP} = & -1 & +1 \end{array}$$

The Fermi model

(Fermi 1933, further developed by Feynman, Gell-Mann and others after 1958)

Lagrangian for “current–current interaction” of four fermions:

$$\mathcal{L}_{\mathrm{Fermi}}(x) = -2\sqrt{2} G_\mu J_\rho^\dagger(x) J^\rho(x), \qquad G_\mu = 1.16639 \times 10^{-5}\,\mathrm{GeV}^{-2}$$

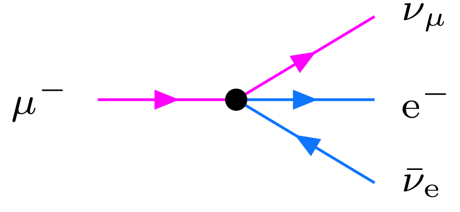
Fermi constant

with $J_\rho(x) = J_\rho^{\mathrm{lep}}(x) + J_\rho^{\mathrm{had}}(x)$ charged weak current

- leptonic current J_ρ^{lep}:

$$J_\rho^{\mathrm{lep}} = \overline{\psi_{\nu_\mathrm{e}}} \gamma_\rho \omega_- \psi_\mathrm{e} + \overline{\psi_{\nu_\mu}} \gamma_\rho \omega_- \psi_\mu, \qquad \omega_\pm = \frac{1}{2}(1 \pm \gamma_5) = \text{chirality projectors}$$

  - ◇ only left-handed fermions $(\omega_- \psi)$, right-handed anti-fermions $(\overline{\psi} \omega_+)$ feel (charged-current) weak interactions ⇒ maximal P-violation
  - ◇ doublet structure: $\begin{pmatrix} \nu_\mathrm{e} \\ \mathrm{e}^- \end{pmatrix}$, $\begin{pmatrix} \nu_\mu \\ \mu^- \end{pmatrix}$, later completed by $\begin{pmatrix} \nu_\tau \\ \tau^- \end{pmatrix}$
  - ◇ $(J^{\mathrm{lep},\rho})^\dagger J_\rho^{\mathrm{lep}}$ induces muon decay:



The Fermi model (cont.)

- Hadronic current J_ρ^{had}:

  formulated in terms of quark fields

  relevant quarks for energies $\lesssim 1\,\text{GeV}$: $\quad \text{u}, \text{d}, \text{s}, \text{c}$

  simple doublet structure $\begin{pmatrix} \text{u} \\ \text{d} \end{pmatrix}, \begin{pmatrix} \text{c} \\ \text{s} \end{pmatrix}$ in conflict with experiment

  e.g. observed process $\underbrace{K^+}_{\text{u}\bar{\text{s}}\text{ pair in quark model}} \to \mu^+ \nu_\mu$ would not be allowed

  solution (Cabibbo 1963):

  u-c-mixing and d-s-mixing in weak interaction

  $\hookrightarrow$ doublets $\begin{pmatrix} \text{u} \\ \text{d}' \end{pmatrix}, \begin{pmatrix} \text{c} \\ \text{s}' \end{pmatrix}$ with $\begin{pmatrix} \text{d}' \\ \text{s}' \end{pmatrix} = U_\text{C} \begin{pmatrix} \text{d} \\ \text{s} \end{pmatrix}$

  orthogonal Cabibbo matrix $\quad U_\text{C} = \begin{pmatrix} \cos\theta_\text{C} & \sin\theta_\text{C} \\ -\sin\theta_\text{C} & \cos\theta_\text{C} \end{pmatrix}$

  empirical result: $\quad \theta_\text{C} \approx 13°$

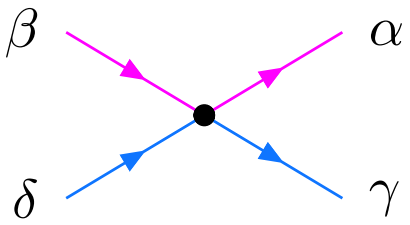
$$J_\rho^{\text{had}} = \overline{\psi_\text{u}} \gamma_\rho \omega_- \psi_{\text{d}'} + \overline{\psi_\text{c}} \gamma_\rho \omega_- \psi_{\text{s}'}$$

Remarks on the Fermi model

- Universality of weak interaction:
  universal coupling G_μ for all transitions and $U_{\mathrm{C}}^\dagger U_{\mathrm{C}} = \mathbf{1}$
- vector and axial-vector interaction sufficient
  to describe low-energy experiments ($E \lesssim 1\,\mathrm{GeV}$)
  no other couplings like (pseudo-)scalar couplings necessary
  $[(\overline{\psi}\psi)(\overline{\psi}\psi), (\overline{\psi}\psi)(\overline{\psi}\gamma_5\psi), \ldots]$
- problems:
  - ◇ cross sections for $\nu_\mu \mathrm{e} \to \nu_{\mathrm{e}}\mu$, etc., grow for energy $E \to \infty$ as E^2

    $\hookrightarrow$ unitarity violation !
  - ◇ no consistent evaluation of higher perturbative orders possible
    (no cancellation of UV divergences)

    $\hookrightarrow$ non-renormalizability !
- elementary interaction:



$$= -\mathrm{i}2\sqrt{2}G_\mu(\gamma_\mu\omega_-)_{\alpha\beta}(\gamma^\mu\omega_-)_{\gamma\delta}$$

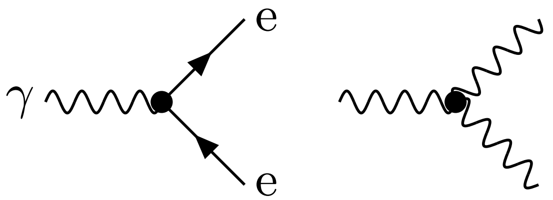
Feynman rules

Lagrangian is equivalent to a set of Feynman rules

propagators for free fields

etc.

vertices = elementary interactions

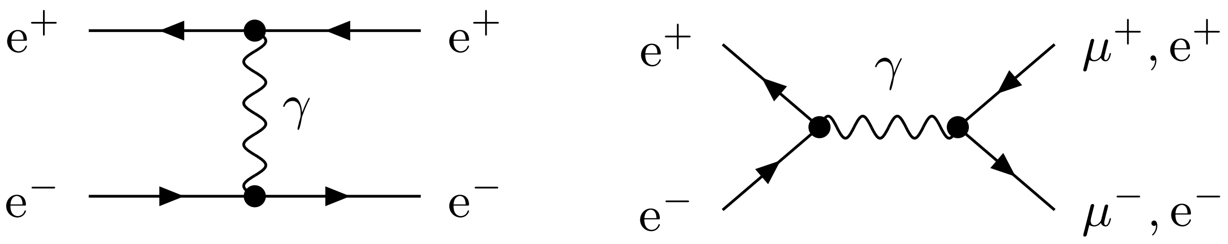


etc.

Feynman diagrams

- provide an exact, clear formulation of perturbation theory
- describe interactions intuitively by scattering processes of free particles

examples for electromagnetic interaction:



transition amplitude (S-matrix element)
$\langle f|S|i\rangle \; = \; \Sigma$ of all Feynman graphs for $|i\rangle \rightarrow |f\rangle$

Intermediate-vector-boson (IVB) model

Idea: “resolution” of four-fermion interaction by vector-boson exchange

Lagrangian:

$$\mathcal{L}_{\mathrm{IVB}} = \mathcal{L}_{0,\mathrm{ferm}} + \mathcal{L}_{0,\mathrm{W}} + \mathcal{L}_{\mathrm{int}}, \qquad \mathcal{L}_{0,\mathrm{ferm}} = \sum_f \overline{\psi_f}(\mathrm{i}\partial\!\!\!/ - m_f)\psi_f$$

$$\mathcal{L}_{0,\mathrm{W}} = -\frac{1}{2}(\partial_\mu W_\nu^+ - \partial_\nu W_\mu^+)(\partial^\mu W^{-,\nu} - \partial^\nu W^{-,\mu}) + M_\mathrm{W}^2 W_\mu^+ W^{-,\mu}$$

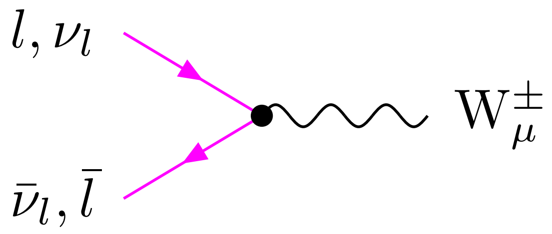
$\mathrm{W}^\pm$ are vector bosons with electric charge $\pm e$ and mass M_W.

W-propagator: $$G_{\mu\nu}^{\mathrm{WW}}(k) = \frac{-\mathrm{i}}{k^2 - M_\mathrm{W}^2}\left(g_{\mu\nu} - \frac{k_\mu k_\nu}{M_\mathrm{W}^2}\right), \qquad k = \text{momentum}$$

interaction Lagrangian: $$\mathcal{L}_{\mathrm{int}} = \frac{g_\mathrm{W}}{\sqrt{2}}\left(J^\rho W_\rho^+ + J^{\rho\dagger} W_\rho^-\right)$$

J^ρ = charged weak current as in Fermi model

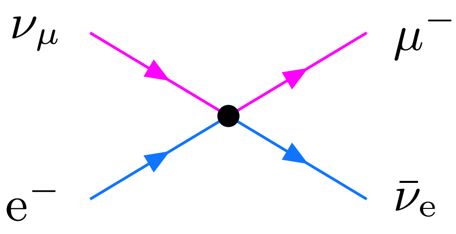
elementary interaction:



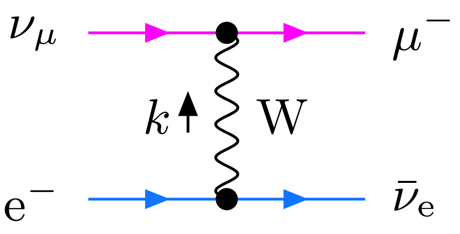
$$= \mathrm{i}\frac{g_\mathrm{W}}{\sqrt{2}}\gamma_\mu \omega_-$$

Four-fermion interaction in process $\nu_\mu e^- \longrightarrow \mu^- \nu_e$

Fermi model:



IVB model:



$$-\mathrm{i}2\sqrt{2}G_\mu\, g_{\rho\sigma}$$

$$\times\left[\bar u_{\mu^-}\gamma^\rho\omega_- u_{\nu_\mu}\right]\,\left[\bar u_{\nu_e}\gamma^\sigma\omega_- u_{e^-}\right]$$

$$\tfrac{\mathrm{i}}{2}g_W^2\frac{1}{k^2-M_W^2}\left(g_{\rho\sigma}-\frac{k_\rho k_\sigma}{M_W^2}\right)$$

$$\times\left[\bar u_{\mu^-}\gamma^\rho\omega_- u_{\nu_\mu}\right]\,\left[\bar u_{\nu_e}\gamma^\sigma\omega_- u_{e^-}\right]$$

identification for $|k| \ll M_W \quad\Rightarrow\quad 2\sqrt{2}G_\mu \;=\; g_W^2/(2M_W^2)$

consequences for the high-energy behaviour:

- k^ρ terms: $\quad \bar u_{\nu_e}\not{k}\omega_- u_{e^-} = \bar u_{\nu_e}(\not{p}_e - \not{p}_{\nu_e})\omega_- u_{e^-} = m_e\bar u_{\nu_e}\omega_+ u_{e^-}$

  $\hookrightarrow$ no extra factors of scattering energy E

- propagator $1/(k^2-M_W^2) \;\sim\; 1/E^2$ for $|k| \sim E \gg M_W$

  $\hookrightarrow$ damping of amplitude in high-energy limit by factor $1/E^2$

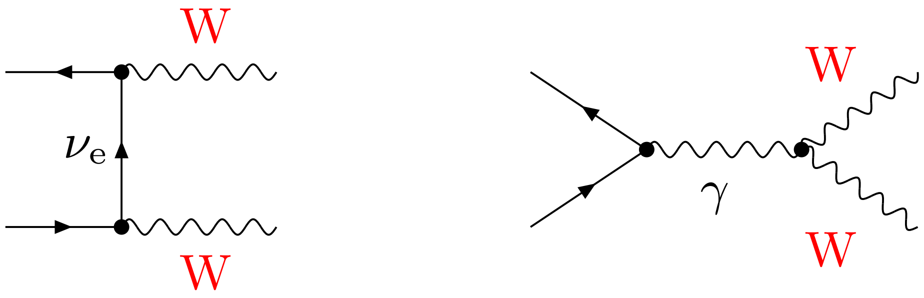
$\Rightarrow$ cross section $\underset{E\to\infty}{\widetilde{\quad}}$ const$/E^2$ $\qquad \Rightarrow$ no unitarity violation !

Comments on the IVB model

- Formal similarity with QED interaction: $J^\rho W^+_\rho$ + h.c. $\longleftrightarrow$ $j^\rho_{\mathrm{elmg.}} A_\rho$
- intermediate vector bosons can be produced, e.g.

$$\underbrace{\mathrm{u\bar{d}}}_{\text{in } \mathrm{p\bar{p}} \text{ collision}} \longrightarrow \underbrace{\mathrm{W}^+ \to f\bar{f}'}_{\mathrm{W}^\pm \text{ unstable}} \qquad \text{(discovery 1983 at CERN)}$$

- problems:
  - ◇ unitarity violations in cross sections with longitudinal W bosons, e.g.



  - ◇ non-renormalizability
    (no consistent treatment of higher perturbative orders)

$\hookrightarrow$ solution by spontaneously broken gauge theories !

Basic principles of the SM

Electroweak Standard Model (SM) constructed in analogy to QED

	QED	SM
matter fields	$e^{\pm}$	leptons, quarks
global symmetry	$U(1)_{em}$	$SU(2)_I \times U(1)_Y$
$\downarrow$ "gauging of the symmetry" global $\rightarrow$ local symmetry		
introduction of gauge bosons and interactions	γ	$\gamma, Z, W^{\pm}$

differences to QED

- non-abelian gauge symmetry $\Rightarrow$ gauge-boson self-interactions
- spontaneous symmetry breaking $SU(2)_I \times U(1)_Y \rightarrow U(1)_{em}$
  $\Rightarrow$ massive gauge bosons $Z, W^{\pm}$ and Higgs boson H (spin 0)

$\Rightarrow$ unified description of electromagnetic and weak interaction
Glashow, Salam, Weinberg 1967–1970:
Standard Model of electroweak interaction (GSW model)

The Principle of local gauge invariance

QED as U(1) gauge theory:

free Lagrangian $\mathcal{L}_{0,\mathrm{ferm}} = \overline{\psi_f}(\mathrm{i}\partial\!\!\!/ - m_f)\psi_f$ invariant under global U(1) symmetry:

$$\psi_f \to \psi'_f = \exp\{-\mathrm{i}Q_f e\theta\}\psi_f, \quad \overline{\psi_f} \to \overline{\psi'_f} = \overline{\psi_f}\exp\{+\mathrm{i}Q_f e\theta\}$$

with space-time-independent group parameter θ

“gauging the symmetry”: demand local symmetry, $\theta \to \theta(x)$

to achieve local symmetry, extend theory by “minimal substitution”:

$$\partial^\mu \to D^\mu = \partial^\mu + \mathrm{i}Q_f e A^\mu(x)$$ = “covariant derivative”,

$A^\mu(x)$ = spin-1 gauge field (photon).

Transformation property of photon $A_\mu(x) \to A'_\mu(x) = A_\mu(x) + \partial_\mu\theta(x)$ ensures

- $D_\mu\psi_f \to (D_\mu\psi_f)' = D'_\mu\psi'_f = \exp\{-\mathrm{i}Q_f e\theta\}(D_\mu\psi_f)$
- gauge invariance of field-strength tensor $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$

gauge-invariant Lagrangian of QED:

$$\mathcal{L}_{\mathrm{QED}} = \underbrace{\overline{\psi_f}(\mathrm{i}\partial\!\!\!/ - Q_f e A\!\!\!/ - m_f)\psi_f}_{\text{fermion part}} - \underbrace{\frac{1}{4}F_{\mu\nu}F^{\mu\nu}}_{\text{gauge part}}$$

Non-Abelian gauge theory (Yang–Mills theory)

Starting point: Lagrangian $\mathcal{L}_\Phi(\Phi, \partial_\mu \Phi)$ of free or self-interacting fields
with “internal symmetry”:

- $\Phi = \begin{pmatrix} \phi_1 \\ \vdots \\ \phi_n \end{pmatrix}$ = multiplet of a compact Lie group G:

  $\Phi \to \Phi' = U(\theta)\Phi, \quad U(\theta) = \exp\{-\mathrm{i} g T^a \theta^a\}$ = unitary,

  T^a = (hermitian) group generators, $\quad a = 1, \ldots, N, \quad N$ = dimension of group

  properties of T^a: $\quad [T^a, T^b] = \mathrm{i} f^{abc} T^c, \quad \mathrm{Tr}\, T^a T^b = \frac{1}{2}\delta^{ab}$

  f^{abc} structure constants of G

- $\mathcal{L}_\Phi$ is invariant under G: $\quad \mathcal{L}_\Phi(\Phi, \partial_\mu \Phi) = \mathcal{L}_\Phi(\Phi', \partial_\mu \Phi')$

examples:

self-interacting (complex) boson multiplet	free fermion multiplet
$\mathcal{L}_\Phi = (\partial_\mu \Phi)^\dagger (\partial^\mu \Phi) - m^2 \Phi^\dagger \Phi + \lambda(\Phi^\dagger \Phi)^2$	$\mathcal{L}_\Psi = \overline{\Psi} \mathrm{i} \partial\!\!\!/ \Psi - m \overline{\Psi} \Psi$
(m = common mass, λ = coupling strength)	

Non-Abelian gauge theory (Yang–Mills theory) (cont.)

Gauging the symmetry by minimal substitution:

$$\mathcal{L}_\Phi(\Phi, \partial_\mu \Phi) \to \mathcal{L}_\Phi(\Phi, D_\mu \Phi) \quad \text{with} \;\; D_\mu = \partial_\mu + \mathrm{i} g T^a A^a_\mu(x)$$

g = gauge coupling, T^a = generator of G in Φ representation
$A^a_\mu(x)$ = gauge fields, $a = 1, \dots, N$

transformation property of gauge fields:

$\mathcal{L}_\Phi(\Phi, D_\mu\Phi)$ local invariant if $\quad D_\mu \Phi \;\to\; (D_\mu\Phi)' = D'_\mu \Phi' = U(\theta)(D_\mu\Phi)$

$$\Rightarrow \; T^a A'^a_\mu \;=\; U T^a A^a_\mu U^\dagger - \frac{\mathrm{i}}{g} U(\partial_\mu U^\dagger), \qquad A^a_\mu A^{a,\mu} = \text{not gauge invariant}$$

infinitesimal form: $\quad \delta A^a_\mu = g f^{abc} \delta\theta^b A^c_\mu + \partial_\mu \delta\theta^a$

covariant definition of field strength: $\quad [D_\mu, D_\nu] = \mathrm{i} g T^a F^a_{\mu\nu}$

$$\Rightarrow \; T^a F^a_{\mu\nu} \;\to\; T^a F'^a_{\mu\nu} = U T^a F^a_{\mu\nu} U^\dagger, \qquad F^a_{\mu\nu} F^{a,\mu\nu} = \text{gauge invariant}$$

explicit form: $\quad F^a_{\mu\nu} = \partial_\mu A^a_\nu - \partial_\nu A^a_\mu - g f^{abc} A^b_\mu A^c_\nu$

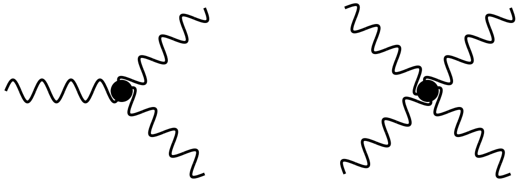
Yang–Mills Lagrangian for gauge and matter fields:

$$\mathcal{L}_{\mathrm{YM}} = -\frac{1}{4} F^a_{\mu\nu} F^{a,\mu\nu} + \mathcal{L}_\Phi(\Phi, D_\mu\Phi)$$

Non-Abelian gauge theory (Yang–Mills theory) (cont.)

Remarks

- Lagrangian contains terms of order $(\partial A)A^2$, A^4 in F^2 part
  $\hookrightarrow$ cubic and quartic gauge-boson self-interactions



- gauge coupling determines gauge-boson–matter and gauge-boson self-interaction $\rightarrow$ unification of interactions
- mass term $M^2(A^a_\mu A^{a,\mu})$ for gauge bosons forbidden by gauge invariance
  $\hookrightarrow$ gauge bosons of unbroken Yang–Mills theory are massless
- non-abelian charges are quantized owing to $[T^a, T^b] = \mathrm{i} f^{abc} T^c$
  (abelian charges are arbitrary)
- $G = \mathrm{SU}(3)$ and fermion triplets $\Rightarrow$ Lagrangian of QCD

$$\mathcal{L}_\Psi = \overline{\Psi} \mathrm{i} \gamma^\mu D_\mu \Psi - m \overline{\Psi} \Psi$$

Choice of the gauge group for electroweak interaction

- Why unification of weak and elmg. interaction ?
  - ◇ similarity: spin-1 fields couple to matter currents formed by spin-$\frac{1}{2}$ fields
  - ◇ elmg. coupling of charged $\mathrm{W}^{\pm}$ bosons
  - ◇ unitarity of theory with elmg. charged massive gauge bosons requires unification

- minimal choice of gauge group: $\mathrm{SU}(2)_{\mathrm{I}} \times \mathrm{U}(1)_{\mathrm{Y}}$
  - ◇ $\mathrm{SU}(2)_{\mathrm{I}}$ $\rightarrow$ weak isospin group with gauge bosons $\mathrm{W}^{+}, \mathrm{W}^{-}, \mathrm{W}^{0}$
    generators I^{a}_{W}, $a = 1, 2, 3$, gauge coupling g_2
  - ◇ $\mathrm{U}(1)_{\mathrm{Y}}$ $\rightarrow$ weak hypercharge group with gauge boson B
    generator Y_{W}, gauge coupling g_1

  W^{0} and B carry identical quantum numbers

  $\hookrightarrow$ two neutral gauge bosons γ, Z as mixed states

  experiment: 1973 discovery of neutral weak currents at CERN
  1983 discovery of Z boson at CERN

Fermion multiplet structure

Distinguish between left-/right-handed parts of fermions: $\psi^{\mathrm{L}} = \omega_{-}\psi$, $\psi^{\mathrm{R}} = \omega_{+}\psi$

- ψ^{L} couple to $\mathrm{W}^{\pm}$ → group ψ^{L} into SU(2)$_{\mathrm{I}}$ doublets, weak isospin $I^a_{\mathrm{w}} = \frac{\sigma^a}{2}$
- ψ^{R} do not couple to $\mathrm{W}^{\pm}$ → ψ^{R} are SU(2)$_{\mathrm{I}}$ singlets, weak isospin $I^a_{\mathrm{w}} = 0$
- $\psi^{\mathrm{L/R}}$ couple to γ in the same way

↪ adjust coupling to $U(1)_{\mathrm{Y}}$ (i.e. fix weak hypercharges $Y^{\mathrm{L/R}}_{\mathrm{w}}$ for $\psi^{\mathrm{L/R}}$)
such that elmg. coupling results: $\mathcal{L}_{\mathrm{int,QED}} = -\sum_f Q_f e \overline{\psi_f} \not{A} \psi_f$

fermion content of the SM:
(ignoring possible right-handed neutrinos)

					T^3_{I}	Q
leptons:	$\Psi^{\mathrm{L}}_{L,i} =$	$\begin{pmatrix}\nu^{\mathrm{L}}_{\mathrm{e}} \\ \mathrm{e}^{\mathrm{L}}\end{pmatrix}$,	$\begin{pmatrix}\nu^{\mathrm{L}}_{\mu} \\ \mu^{\mathrm{L}}\end{pmatrix}$,	$\begin{pmatrix}\nu^{\mathrm{L}}_{\tau} \\ \tau^{\mathrm{L}}\end{pmatrix}$,	$+\frac{1}{2}$ $-\frac{1}{2}$	0 -1
	$\psi^{\mathrm{R}}_{l,i} =$	e^{R},	μ^{R},	τ^{R},	0	-1
quarks: (each quark exists in 3 colours!)	$\Psi^{\mathrm{L}}_{Q,i} =$	$\begin{pmatrix}\mathrm{u}^{\mathrm{L}} \\ \mathrm{d}^{\mathrm{L}}\end{pmatrix}$,	$\begin{pmatrix}\mathrm{c}^{\mathrm{L}} \\ \mathrm{s}^{\mathrm{L}}\end{pmatrix}$,	$\begin{pmatrix}\mathrm{t}^{\mathrm{L}} \\ \mathrm{b}^{\mathrm{L}}\end{pmatrix}$,	$+\frac{1}{2}$ $-\frac{1}{2}$	$+\frac{2}{3}$ $-\frac{1}{3}$
	$\psi^{\mathrm{R}}_{u,i} =$	u^{R},	c^{R},	t^{R},	0	$+\frac{2}{3}$
	$\psi^{\mathrm{R}}_{d,i} =$	d^{R},	s^{R},	b^{R},	0	$-\frac{1}{3}$

Fermion Lagrangian and minimal substitution

Free Lagrangian of (still massless) fermions:

$$\mathcal{L}_{0,\mathrm{ferm}}=\sum_f \mathrm{i}\overline{\psi_f}\partial\!\!\!/\psi_f$$

$$=\sum_i(\mathrm{i}\overline{\Psi^{\mathrm{L}}_{L,i}}\partial\!\!\!/\Psi^{\mathrm{L}}_{L,i}+\mathrm{i}\overline{\Psi^{\mathrm{L}}_{Q,i}}\partial\!\!\!/\Psi^{\mathrm{L}}_{Q,i}+\mathrm{i}\overline{\psi^{\mathrm{R}}_{l,i}}\partial\!\!\!/\psi^{\mathrm{R}}_{l,i}+\mathrm{i}\overline{\psi^{\mathrm{R}}_{u,i}}\partial\!\!\!/\psi^{\mathrm{R}}_{u,i}+\mathrm{i}\overline{\psi^{\mathrm{R}}_{d,i}}\partial\!\!\!/\psi^{\mathrm{R}}_{d,i})$$

minimal substitution: $\partial_\mu \;\rightarrow\; D_\mu$

$$D_\mu=\partial_\mu-\mathrm{i}g_2I^a_{\mathrm{w}}W^a_\mu+\mathrm{i}g_1\tfrac{1}{2}Y_{\mathrm{w}}B_\mu=D^{\mathrm{L}}_\mu\omega_-+D^{\mathrm{R}}_\mu\omega_+,$$

$$D^{\mathrm{L}}_\mu=\partial_\mu-\frac{\mathrm{i}g_2}{\sqrt{2}}\begin{pmatrix}0 & W^+_\mu\\ W^-_\mu & 0\end{pmatrix}+\frac{\mathrm{i}}{2}\begin{pmatrix}-g_2W^3_\mu+g_1Y^{\mathrm{L}}_{\mathrm{w}}B_\mu & 0\\ 0 & g_2W^3_\mu+g_1Y^{\mathrm{L}}_{\mathrm{w}}B_\mu\end{pmatrix}$$

$$D^{\mathrm{R}}_\mu=\partial_\mu+\mathrm{i}g_1\tfrac{1}{2}Y^{\mathrm{R}}_{\mathrm{w}}B_\mu$$

charge eigenstates: $W^\pm_\mu=\frac{1}{\sqrt{2}}(W^1_\mu\mp\mathrm{i}W^2_\mu),\quad W^3_\mu,\quad B_\mu$

$$QW^\pm_\mu=I^3_{\mathrm{w}}W^\pm_\mu=\pm W^\pm_\mu,\qquad QW^3_\mu=QB_\mu=0$$

with $Q=I^3_{\mathrm{w}}+Y_{\mathrm{w}}/2$ (see below)

Photon identification

"Weinberg rotation": $\begin{pmatrix} Z_\mu \\ A_\mu \end{pmatrix} = \begin{pmatrix} c_\mathrm{w} & s_\mathrm{w} \\ -s_\mathrm{w} & c_\mathrm{w} \end{pmatrix} \begin{pmatrix} W^3_\mu \\ B_\mu \end{pmatrix},$

$c_\mathrm{w} = \cos\theta_\mathrm{w}, \; s_\mathrm{w} = \sin\theta_\mathrm{w},$ $\quad \theta_\mathrm{w} =$ Weinberg angle = electroweak mixing angle

$$D^\mathrm{L}_\mu\big|_{A_\mu} = \frac{\mathrm{i}}{2} A_\mu \begin{pmatrix} g_2 s_\mathrm{w} + g_1 c_\mathrm{w} Y^\mathrm{L}_\mathrm{w} & 0 \\ 0 & -g_2 s_\mathrm{w} + g_1 c_\mathrm{w} Y^\mathrm{L}_\mathrm{w} \end{pmatrix} \overset{!}{=} \mathrm{i} e A_\mu \begin{pmatrix} Q_1 & 0 \\ 0 & Q_2 \end{pmatrix}$$

$$D^\mathrm{R}_\mu\big|_{A_\mu} = \frac{\mathrm{i}}{2} A_\mu g_1 c_\mathrm{w} Y^\mathrm{R}_\mathrm{w} \overset{!}{=} \mathrm{i} e A_\mu Q$$

- charge difference in doublet $Q_1 - Q_2 = 1$ $\quad \to \quad g_2 = \dfrac{e}{s_\mathrm{w}}$
- normalize $Y^{\mathrm{L/R}}_\mathrm{w}$ such that $\quad g_1 = \dfrac{e}{c_\mathrm{w}}$

$\hookrightarrow$ Y_w fixed by "Gell-Mann–Nishijima relation": $\quad Q = I^3_\mathrm{w} + \dfrac{Y_\mathrm{w}}{2}$

parameter relations: $e = \dfrac{g_1 g_2}{\sqrt{g_1^2 + g_2^2}}, \qquad \tan\theta_\mathrm{w} = \dfrac{g_1}{g_2}$

$$c_\mathrm{w} = \frac{g_2}{\sqrt{g_1^2 + g_2^2}}, \qquad s_\mathrm{w} = \frac{g_1}{\sqrt{g_1^2 + g_2^2}}$$

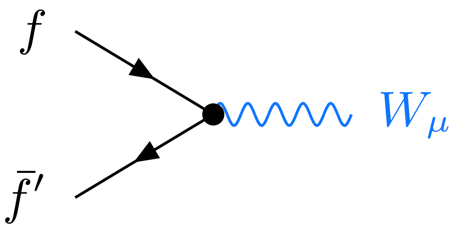
Fermion–gauge-boson interaction:

$$\mathcal{L}_{\mathrm{ferm,YM}} = \sum_F \left(\frac{e}{\sqrt{2}s_{\mathrm{w}}} \overline{\Psi_F^{\mathrm{L}}} \begin{pmatrix} 0 & \not{W}^+ \\ \not{W}^- & 0 \end{pmatrix} \Psi_F^{\mathrm{L}} + \frac{e}{2c_{\mathrm{w}}s_{\mathrm{w}}} \overline{\Psi_F^{\mathrm{L}}} \sigma^3 \not{Z} \Psi_F^{\mathrm{L}} \right)$$

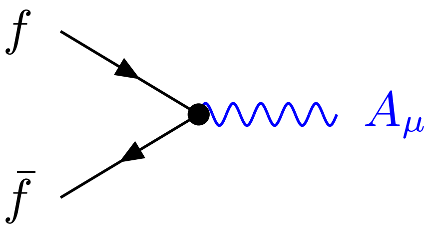
$$- \sum_f \left(e\frac{s_{\mathrm{w}}}{c_{\mathrm{w}}} Q_f \overline{\psi_f} \not{Z} \psi_f + e Q_f \overline{\psi_f} \not{A} \psi_f \right)$$

(f=all fermions, F= all doublets)

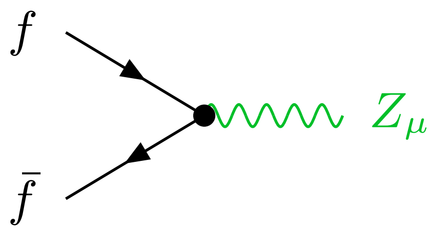
Feynman rules:



$$\frac{\mathrm{i}e}{\sqrt{2}s_{\mathrm{w}}}\gamma_\mu\omega_-$$



$$-\mathrm{i}Q_f e\gamma_\mu$$



$$\mathrm{i}e\gamma_\mu(g_f^+\omega_+ + g_f^-\omega_-) = \mathrm{i}e\gamma_\mu(v_f - a_f\gamma_5)$$

with $g_f^+ = -\dfrac{s_{\mathrm{w}}}{c_{\mathrm{w}}}Q_f, \quad g_f^- = -\dfrac{s_{\mathrm{w}}}{c_{\mathrm{w}}}Q_f + \dfrac{I_{\mathrm{w},f}^3}{c_{\mathrm{w}}s_{\mathrm{w}}}$

$$v_f = -\frac{s_{\mathrm{w}}}{c_{\mathrm{w}}}Q_f + \frac{I_{\mathrm{w},f}^3}{2c_{\mathrm{w}}s_{\mathrm{w}}}, \quad a_f = \frac{I_{\mathrm{w},f}^3}{2c_{\mathrm{w}}s_{\mathrm{w}}}$$

Gauge-boson sector

Yang–Mills Lagrangian for gauge fields:

$$\mathcal{L}_{\mathrm{YM}} = -\frac{1}{4} W^a_{\mu\nu} W^{a,\mu\nu} - \frac{1}{4} B_{\mu\nu} B^{\mu\nu}$$

field-strength tensors:

$$W^a_{\mu\nu} \;=\; \partial_\mu W^a_\nu - \partial_\nu W^a_\mu + g_2 \epsilon^{abc} W^b_\mu W^c_\nu, \qquad B_{\mu\nu} \;=\; \partial_\mu B_\nu - \partial_\nu B_\mu$$

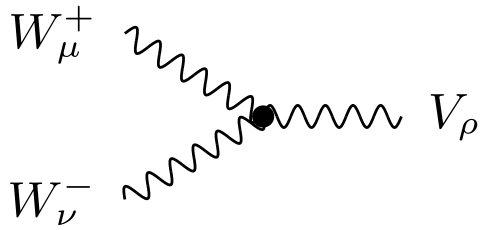
Yang–Mills Lagrangian in terms of “physical” fields:

$$\begin{aligned}\mathcal{L}_{\mathrm{YM}} = &-\frac{1}{4}\left|\partial_\mu A_\nu - \partial_\nu A_\mu - \mathrm{i}e(W^-_\mu W^+_\nu - W^-_\nu W^+_\mu)\right|^2 \\ &- \frac{1}{4}\left|\partial_\mu Z_\nu - \partial_\nu Z_\mu + \mathrm{i}e\frac{c_{\mathrm{w}}}{s_{\mathrm{w}}}(W^-_\mu W^+_\nu - W^-_\nu W^+_\mu)\right|^2 \\ &- \frac{1}{2}\left|\partial_\mu W^+_\nu - \partial_\nu W^+_\mu - \mathrm{i}e(W^+_\mu A_\nu - W^+_\nu A_\mu) + \mathrm{i}e\frac{c_{\mathrm{w}}}{s_{\mathrm{w}}}(W^+_\mu Z_\nu - W^+_\nu Z_\mu)\right|^2\end{aligned}$$

$\hookrightarrow$ triple gauge-boson couplings AW^+W^-, ZW^+W^-
quartic gauge-boson couplings AAW^+W^-, AZW^+W^-, ZZW^+W^-, $W^+W^-W^+W^-$

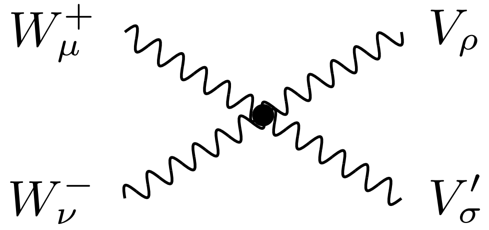
Feynman rules for gauge-boson self-interactions

(fields and momenta incoming)



$$\mathrm{i}e\,C_{WWV}\Big[g_{\mu\nu}(k_+-k_-)_\rho+g_{\nu\rho}(k_--k_V)_\mu + g_{\rho\mu}(k_V-k_+)_\nu\Big]$$

with $C_{WW\gamma}=1,\quad C_{WWZ}=-\dfrac{c_\mathrm{w}}{s_\mathrm{w}}$



$$\mathrm{i}e^2C_{WWVV'}\Big[2g_{\mu\nu}g_{\rho\sigma}-g_{\mu\rho}g_{\sigma\nu}-g_{\mu\sigma}g_{\nu\rho}\Big]$$

with $C_{WW\gamma\gamma}=-1,\qquad C_{WW\gamma Z}=\dfrac{c_\mathrm{w}}{s_\mathrm{w}}$

$C_{WWZZ}=-\dfrac{c_\mathrm{w}^2}{s_\mathrm{w}^2},\quad C_{WWWW}=\dfrac{1}{s_\mathrm{w}^2}$

Introduction of gauge-boson masses

Consistency of theory (unitarity, renormalizability)
$\Leftarrow$ gauge symmetry of Lagrangian

explicit gauge-boson mass terms violate gauge invariance

solution: spontaneous symmetry breaking (hidden symmetry)

- invariant Lagrangian
- non-invariant ground state

Standard Model: Higgs mechanism

introduce scalar field with non-vanishing vacuum expectation value (vev) that couples to gauge bosons

(non-zero vev of fields with spin violates Lorentz invariance)

idea: spontaneous breakdown $\mathrm{SU(2)_I \times U(1)_Y \;\to\; U(1)_{elmg}}$

$\hookrightarrow$ masses for $\mathrm{W}^{\pm}$ and Z bosons, but γ remains massless

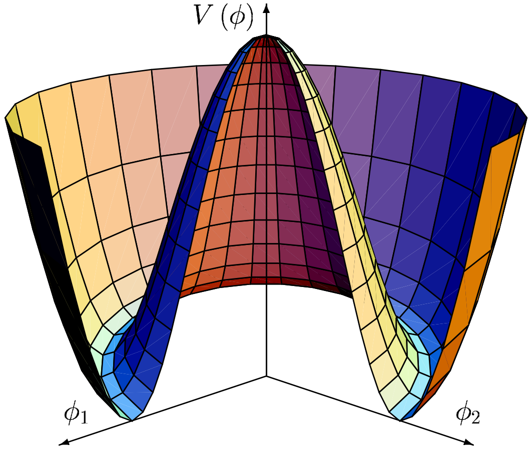
Note: choice of scalar extension of massless model involves freedom

Higgs sector and spontaneous symmetry breaking in SM

GSW model:

minimal scalar sector with complex scalar doublet $\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$, $Y_\mathrm{w}^\Phi = 1$

scalar self-interaction via Higgs potential:



$$V(\Phi) = -\mu^2\Phi^\dagger\Phi + \frac{\lambda}{4}(\Phi^\dagger\Phi)^2, \quad \mu^2, \lambda > 0,$$

$$= \mathrm{SU(2)_I}\times\mathrm{U(1)_Y} \text{ symmetric}$$

$$V(\Phi) = \text{minimal for} \quad |\Phi| = \sqrt{\frac{2\mu^2}{\lambda}} \equiv \frac{v}{\sqrt{2}} > 0$$

ground state Φ_0 (= vacuum expectation value of Φ) not unique

choice $\Phi_0 = \begin{pmatrix} 0 \\ \frac{v}{\sqrt{2}} \end{pmatrix}$ not gauge invariant $\Rightarrow$ spontaneous symmetry breaking

elmg. gauge invariance unbroken, since $\quad Q\Phi_0 = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \Phi_0 = 0$

field excitations in Φ:

$$\Phi(x) = \begin{pmatrix} \phi^+(x) \\ \frac{1}{\sqrt{2}}\left(v + H(x) + \mathrm{i}\chi(x)\right) \end{pmatrix}$$

Higgs mechanism

spontaneous breaking of a global symmetry

Goldstone theorem: for each spontaneously broken symmetry exists
one massless scalar boson
= Goldstone boson (excitation along minimum of potential)

spontaneous breaking of a local symmetry (gauge symmetry)

Higgs mechanism: degrees of freedom of Goldstone bosons are transmuted into longitudinal degrees of freedom of massless gauge bosons

(would-be) Goldstone bosons are unphysical degrees of freedom:
gauge-dependent masses, absent in unitary gauge

Standard Model

need three (real) longitudinal degrees of freedom for massive $\mathrm{Z}, \mathrm{W}^{\pm}$
$\hookrightarrow$ three components of scalar field(s) are transmuted
other components appear as physical scalar fields
complex Higgs doublet (4 d.o.f) $\Rightarrow$ one physical scalar field

Higgs-Lagrangian of the Electroweak Standard Model

Gauge-invariant Lagrangian of Higgs sector: $(\phi^- = (\phi^+)^\dagger)$

$$\mathcal{L}_\mathrm{H} = (D_\mu\Phi)^\dagger(D^\mu\Phi) - V(\Phi) \qquad \text{with } D_\mu = \partial_\mu - \mathrm{i}g_2\frac{\sigma^a}{2}W_\mu^a + \mathrm{i}\frac{g_1}{2}B_\mu$$

$$= (\partial_\mu\phi^+)(\partial^\mu\phi^-) - \frac{\mathrm{i}ev}{2s_\mathrm{w}}(W_\mu^+\partial^\mu\phi^- - W_\mu^-\partial^\mu\phi^+) + \frac{e^2v^2}{4s_\mathrm{w}^2}W_\mu^+W^{-,\mu}$$

$$+ \frac{1}{2}(\partial\chi)^2 + \frac{ev}{2c_\mathrm{w}s_\mathrm{w}}Z_\mu\partial^\mu\chi + \frac{e^2v^2}{4c_\mathrm{w}^2s_\mathrm{w}^2}Z^2 + \frac{1}{2}(\partial H)^2 - \mu^2H^2$$

$+$ (interaction terms)

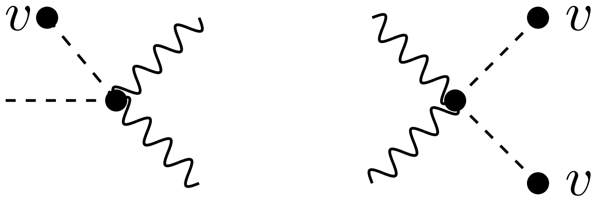
implications:

- gauge-boson masses: $M_\mathrm{W} = \frac{ev}{2s_\mathrm{w}}, \quad M_\mathrm{Z} = \frac{ev}{2c_\mathrm{w}s_\mathrm{w}} = \frac{M_\mathrm{W}}{c_\mathrm{w}}$
- ρ-parameter: $\rho \equiv \frac{M_\mathrm{W}^2}{M_\mathrm{Z}^2c_\mathrm{w}^2} = 1$ (for Higgs doublets!)
- photon remains massless owing to unbroken $\mathrm{U}(1)_\mathrm{em}$ invariance
- physical Higgs boson H: $M_\mathrm{H} = \sqrt{2\mu^2}$ = free parameter
- no mass terms for would-be Goldstone bosons $\phi^\pm$, χ
- mixing terms between gauge fields and would-be Goldstone-boson fields

Higgs-boson interactions

gauge-boson–Higgs-boson couplings

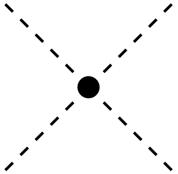
gauge couplings



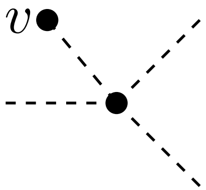
from spontaneous symmetry breaking

$\hookrightarrow$ gauge-boson masses

Higgs-boson self-couplings



$\propto \frac{M_\mathrm{H}^2}{M_\mathrm{W}^2}$ from Higgs potential



$\propto \frac{M_\mathrm{H}^2}{M_\mathrm{W}}$ from spontaneous symmetry breaking

Fermion masses and Yukawa couplings

Ordinary Dirac mass terms $m_f \overline{\psi_f} \psi_f = m_f (\overline{\psi_f^{\rm L}} \psi_f^{\rm R} + \overline{\psi_f^{\rm R}} \psi_f^{\rm L})$ not gauge invariant

$\hookrightarrow$ introduce fermion masses by (gauge-invariant) Yukawa interaction

Lagrangian for Yukawa couplings: (needs Higgs doublets with $Y_{\rm w} = \pm 1$)

$$\mathcal{L}_{\rm Yuk} \;=\; -\overline{\Psi'^{\rm L}_L} G_l \psi'^{\rm R}_l \Phi - \overline{\Psi'^{\rm L}_Q} G_u \psi'^{\rm R}_u \tilde{\Phi} - \overline{\Psi'^{\rm L}_Q} G_d \psi'^{\rm R}_d \Phi + \text{h.c.}$$

- G_l, G_u, G_d = 3×3 matrices in 3-dim. space of generations (ν masses ignored)
- $\tilde{\Phi} = i\sigma^2 \Phi^* = \begin{pmatrix} \phi^{0*} \\ -\phi^- \end{pmatrix}$ = charge conjugate Higgs doublet, $Y_{\rm w}^{\tilde{\Phi}} = -1$

fermion mass terms:

mass terms = bilinear terms in $\mathcal{L}_{\rm Yuk}$, obtained by setting $\Phi \to \Phi_0$:

$$\mathcal{L}_{m_f} \;=\; -\frac{v}{\sqrt{2}} \overline{\psi'^{\rm L}_l} G_l \psi'^{\rm R}_l - \frac{v}{\sqrt{2}} \overline{\psi'^{\rm L}_u} G_u \psi'^{\rm R}_u - \frac{v}{\sqrt{2}} \overline{\psi'^{\rm L}_d} G_d \psi'^{\rm R}_d + \text{h.c.}$$

$\hookrightarrow$ diagonalization by unitary field transformations $\quad (f = l, u, d)$

$\psi_f^{\rm L/R} \equiv U_f^{\rm L/R} \psi_f'^{\rm L/R}$ such that $\frac{v}{\sqrt{2}} U_f^{\rm L} G_f (U_f^{\rm R})^\dagger = {\rm diag}(m_f)$

$\Rightarrow$ standard form: $\mathcal{L}_{m_f} \;=\; -m_f \overline{\psi_f^{\rm L}} \psi_f^{\rm R} + \text{h.c.} \;=\; -m_f \overline{\psi_f} \psi_f$

Quark mixing

- ψ'_f correspond to eigenstates of the gauge interaction
- ψ_f correspond to mass eigenstates,
  for massless neutrinos define $\psi_\nu^{\mathrm{L}} \equiv U_l^{\mathrm{L}} \psi_\nu^{\mathrm{L}}$ → no lepton-flavour change

Yukawa and gauge interactions in terms of mass eigenstates:

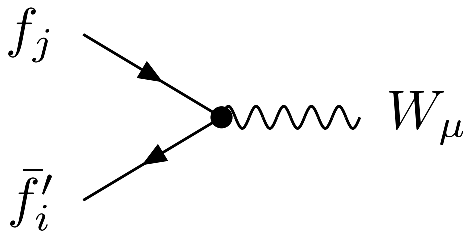
$$\begin{aligned}
\mathcal{L}_{\mathrm{Yuk}} = & -\frac{\sqrt{2}m_l}{v}\left(\phi^+\overline{\psi^{\mathrm{L}}_{\nu_l}}\psi_l^{\mathrm{R}} + \phi^-\overline{\psi_l^{\mathrm{R}}}\psi^{\mathrm{L}}_{\nu_l}\right) + \frac{\sqrt{2}m_u}{v}\left(\phi^+\overline{\psi_u^{\mathrm{R}}}\mathbf{V}\psi_d^{\mathrm{L}} + \phi^-\overline{\psi_d^{\mathrm{L}}}\mathbf{V}^\dagger\psi_u^{\mathrm{R}}\right)\\
& -\frac{\sqrt{2}m_d}{v}\left(\phi^+\overline{\psi_u^{\mathrm{L}}}\mathbf{V}\psi_d^{\mathrm{R}} + \phi^-\overline{\psi_d^{\mathrm{R}}}\mathbf{V}^\dagger\psi_u^{\mathrm{L}}\right) - \frac{m_f}{v}\mathrm{i}\,\mathrm{sgn}(I^3_{\mathrm{w},f})\chi\,\overline{\psi_f}\gamma_5\psi_f\\
& -\frac{m_f}{v}(v+H)\,\overline{\psi_f}\psi_f,
\end{aligned}$$

$$\begin{aligned}
\mathcal{L}_{\mathrm{ferm,YM}} = & \frac{e}{\sqrt{2}s_{\mathrm{w}}}\overline{\Psi^{\mathrm{L}}_L}\begin{pmatrix} 0 & \not{W}^+ \\ \not{W}^- & 0 \end{pmatrix}\psi_L^{\mathrm{L}} + \frac{e}{\sqrt{2}s_{\mathrm{w}}}\overline{\Psi^{\mathrm{L}}_Q}\begin{pmatrix} 0 & \mathbf{V}\not{W}^+ \\ \mathbf{V}^\dagger\not{W}^- & 0 \end{pmatrix}\psi_Q^{\mathrm{L}}\\
& + \frac{e}{2c_{\mathrm{w}}s_{\mathrm{w}}}\overline{\Psi^{\mathrm{L}}_F}\sigma^3\not{Z}\Psi_F^{\mathrm{L}} - e\frac{s_{\mathrm{w}}}{c_{\mathrm{w}}}Q_f\overline{\psi_f}\not{Z}\psi_f - eQ_f\overline{\psi_f}\not{A}\psi_f
\end{aligned}$$

- only charged-current coupling of quarks modified by $\mathbf{V} = U_u^{\mathrm{L}}(U_d^{\mathrm{L}})^\dagger$ = unitary
  (Cabibbo–Kobayashi–Maskawa (CKM) matrix)
- Higgs–fermion coupling strength = m_f/v

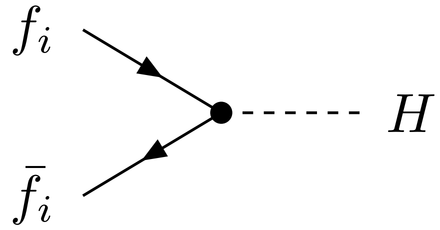
Feynman rules for fermion interactions

quark-mixing matrix $\mathbf{V}_{ij}$ in $\mathrm{W}ff'$ couplings

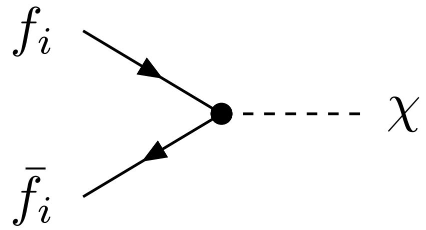


$$\frac{\mathrm{i}e}{\sqrt{2}s_{\mathrm{W}}}\mathbf{V}_{ij}\gamma_\mu\omega_-$$

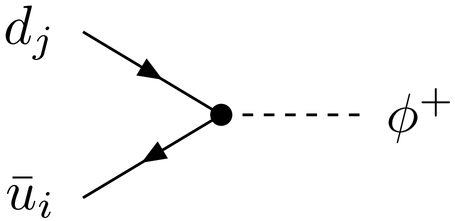
Higgs-boson–fermion couplings (Yukawa couplings)



$$-\frac{\mathrm{i}e}{2s_{\mathrm{W}}}\frac{m_f}{M_{\mathrm{W}}}$$



$$-\frac{e}{2s_{\mathrm{W}}}2I^3_{\mathrm{W},f}\frac{m_f}{M_{\mathrm{W}}}\gamma_5$$



$$-\frac{\mathrm{i}e}{\sqrt{2}s_{\mathrm{W}}}\mathbf{V}_{ij}\left(\frac{m_{u,i}}{M_{\mathrm{W}}}\omega_- - \frac{m_{d,j}}{M_{\mathrm{W}}}\omega_+\right)$$

Features of the CKM mixing

- $\mathbf{V}$ = 3-dim. generalization of Cabibbo matrix U_C
- $\mathbf{V}$ is parametrized by 4 free parameters: 3 real angles, 1 complex phase
  $\hookrightarrow$ complex phase is the only source of CP violation in SM

  counting:

  $$\begin{pmatrix}\text{\#real d.o.f.}\\ \text{in } \mathbf{V}\end{pmatrix} - \begin{pmatrix}\text{\#unitarity}\\ \text{relations}\end{pmatrix} - \begin{pmatrix}\text{\#phase diffs. of}\\ u\text{-type quarks}\end{pmatrix} - \begin{pmatrix}\text{\#phase diffs. of}\\ d\text{-type quarks}\end{pmatrix} - \begin{pmatrix}\text{\#phase diff. between}\\ u\text{- and } d\text{-type quarks}\end{pmatrix}$$

  $$= 18 - 9 - 2 - 2 - 1 = 4$$
- no flavour-changing neutral currents in lowest order,
  flavour-changing suppressed by factors $G_\mu(m_{q_1}^2 - m_{q_2}^2)$ in higher orders
  ("Glashow–Iliopoulos–Maiani mechanism")

Lagrangian of Electroweak Standard Model

Sum of all contributions (classical Lagrangian)

$$\mathcal{L}_{\mathrm{class}} = \mathcal{L}_{\mathrm{ferm,YM}} + \mathcal{L}_{\mathrm{YM}} + \mathcal{L}_{\mathrm{H}} + \mathcal{L}_{\mathrm{Yuk}}$$

contains all possible terms that

- can be build from fields of Standard Model
- are gauge-invariant
- are renormalizable

these restrictions imply perturbative baryon-number conservation and lepton-number conservation

addition of right-handed neutrinos $\nu_{\mathrm{e}}^{\mathrm{R}}$, ν_{μ}^{R}, ν_{τ}^{R}:
$\Rightarrow$ neutrino masses, mixing matrix in lepton sector

without lepton-number violation $\Rightarrow$ analogously as in quark sector

right-handed neutrinos allow for lepton-number violation

Physical parameters of the Electroweak Standard Model

- gauge sector

  $g_1, g_2 \to e = \frac{g_1 g_2}{\sqrt{g_1^2+g_2^2}}, \qquad \cos\theta_{\mathrm{w}} = \frac{g_2}{\sqrt{g_1^2+g_2^2}} = \frac{M_{\mathrm{W}}}{M_{\mathrm{Z}}}$

  elementary charge, weak mixing angle
  (2 parameters)

- Higgs sector

  $\lambda, \mu \to M_{\mathrm{H}} = \sqrt{2}\mu, \qquad M_{\mathrm{W}} = \frac{g_2}{2} v \quad (v = \frac{2\mu}{\sqrt{\lambda}})$

  Higgs-boson mass, W-boson mass
  (2 parameters)

  $M_{\mathrm{Z}} = \frac{\sqrt{g_2^2+g_1^2}}{2} v$ Z-boson mass, weak mixing angle

- flavour sector

  $G_{l,ij}, G_{u,ij}, G_{d,ij} \to m_{f,i} = \frac{1}{\sqrt{2}} \sum_{k,m} U^{\mathrm{L}}_{f,ik} G_{f,km} U^{\mathrm{R}\dagger}_{f,mi} v, \qquad \mathbf{V} = U^{\mathrm{L}}_u U^{\mathrm{L}\dagger}_d$

  fermion masses, quark-mixing matrix
  (9+4 parameters)

  with ν^{R}: 12+8 parameters (neutrino masses, lepton-mixing matrix)

Features of the electroweak Standard Model

- Higgs boson not yet found, particle content verified otherwise

  experimental limit from direct searches: $M_{\mathrm{H}} > 114.1\,\mathrm{GeV}$ 90% CL

  LEPHIGGS '01

  consistency of theory requires: $M_{\mathrm{H}} \lesssim 600\,\mathrm{GeV}$

- No really significant contradictions of GSW model with experiment

- Input parameters:

  $$\alpha = \frac{e^2}{4\pi} \approx 1/137, \;\; M_{\mathrm{W}} \approx 80\,\mathrm{GeV}, \;\; M_{\mathrm{Z}} \approx 91\,\mathrm{GeV}, \;\; M_{\mathrm{H}} \gtrsim 100\,\mathrm{GeV}, \;\; m_f, \;\; V$$

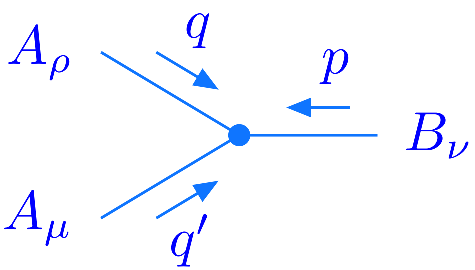
- GSW model = consistent quantum field theory
  - ◇ matrix elements respect unitarity
  - ◇ renormalizability

  ⇒ evaluation of higher perturbative orders possible
  (and phenomenologically necessary !)

Recipe for derivation of Feynman rules

1. Collect all monomials in $\mathrm{i}\mathcal{L}$ that contain a specific selection of fields; these fields form the external legs of a vertex (flow of quantum numbers: incoming)

   e.g.: $-e(\partial_\mu A_\nu)A^\mu B^\nu = -e g_{\rho\nu}(\partial_\mu A^\rho)A^\mu B^\nu$



2. replace all derivatives by $(-\mathrm{i})$ times the incoming momenta of the fields to which they apply
   $\hookrightarrow \; -e(-\mathrm{i}q_\mu)g_{\rho\nu}A^\rho A^\mu B^\nu$
3. sum over all permutations of indices and momenta of identical external fields
   $\hookrightarrow \; \mathrm{i}e(q_\mu g_{\rho\nu} + q'_\rho g_{\mu\nu})A^\rho A^\mu B^\nu$
4. drop all factors of external fields
   $\hookrightarrow \; \mathrm{i}e(q_\mu g_{\rho\nu} + q'_\rho g_{\mu\nu})$

propagator = negative inverse 2-point vertex function

Literature

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